

# Vessiot's Equivalence Method Applied to Linear Partial Differential Operators

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15.5.2008



# Outline

- 1 Motivation
- 2 Vessiot Equivalence Method
  - Introductory Example
  - Natural Bundles
  - Symmetry Groupoids
  - Integrability Conditions
  - Equivalence
- 3 Applications to LPDOs
  - Groupoids  $\Theta_q$  of Gauge Transformations
  - Natural  $\Theta_q$ -Bundles
  - Invariants
- 4 Examples

# Motivation

## Example

Linear partial differential operators (LPDOs) of order 2:

$$L = D_{x^1}D_{x^2} + a D_{x^1} + b D_{x^2} + c$$

under gauge transformations:

$$L \mapsto g^{-1}Lg, \quad g = g(x^1, x^2).$$

Questions:

- Equivalence:  $L, L' \rightsquigarrow g^{-1}Lg = L'?$
- Invariants?

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Questions:

- Equivalence:  $L, L' \rightsquigarrow g^{-1}Lg = L'?$
- Invariants? Laplace:

$$h = c - a_{x^1} - ab, \quad k = c - b_{x^2} - ab.$$

- Generating set of invariants?
- Larger examples?

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# Natural bundles

Let  $X$  be a manifold, coordinates  $(x) = (x^1, \dots, x^n)$ .

- $\text{Diff}_{\text{loc}}(X, X)$ : local diffeomorphisms  $\varphi : X \rightarrow X$ .
- A **natural bundle** is a fibre bundle

$$\pi : \mathcal{F} \rightarrow X : (x, u) \rightarrow (x)$$

such that each  $\varphi \in \text{Diff}_{\text{loc}}(X, X)$  continues to  $\tilde{\varphi} : \mathcal{F} \rightarrow \mathcal{F}$ .

- A section  $\omega : X \rightarrow \mathcal{F} : (x) \mapsto (x, u = \omega(x))$  is called **geometric object**.

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$$\omega \circ \tilde{\varphi} = \gamma$$

- $\varphi$  is a **symmetry** of  $\omega \Leftrightarrow$

$$\omega \circ \tilde{\varphi} = \omega, \quad \Phi_{\omega(y)}(\varphi, \partial_x \varphi, \dots, \varphi_q) = \omega(x).$$

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- $\psi : \mathcal{F} \rightarrow \mathbb{R}$  is an **invariant** if  $\psi \circ \tilde{\varphi} = \psi \quad \forall \varphi \in \text{Diff}_{\text{loc}}(X, X)$ .

# Symmetry Groupoids

- The jet groupoid  $\Pi_q = \Pi_q(X, X)$  has coordinates:

$$(x, y, y_x, \dots, y_q).$$

- Like  $\text{Diff}_{\text{loc}}(X, X)$ ,  $\Pi_q$  acts on  $\mathcal{F}$ .
- Symmetry groupoid  $\mathcal{R}_q(\omega)$  of  $\omega : X \rightarrow \mathcal{F}$  defined by:

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- Prolongation

$$0 \longrightarrow \mathcal{R}_{q+1}(\omega) \longrightarrow \Pi_{q+1} \begin{array}{c} \xrightarrow{j_1(\Phi_\omega)} \\ \xrightarrow{j_1(\omega)} \end{array} J_1(\mathcal{F})$$

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Question: Is  $\mathcal{R}_q^{(1)}(\omega) = \mathcal{R}_q(\omega)$ ?

# Integrability Conditions

## Theorem

$\pi_q^{q+1} : \mathcal{R}_{q+1}(\omega) \rightarrow \mathcal{R}_q(\omega)$  is surjective if and only if there is a section  $c : \mathcal{F} \rightarrow \mathcal{F}_1$ :

- $c$  is equivariant:

$$c(af_q) = c(a)f_q \quad \forall f_q \in \Pi_q, a \in \mathcal{F}$$

- $c$  fulfills the **Vessiot structure equations**:

$$(I \circ j_1)(\omega) = c(\omega).$$

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If in addition the symbol of  $\mathcal{R}_q(\omega)$  is 2-acyclic,  $\mathcal{R}_q(\omega)$  is **integrable**, i. e.:

$$\pi_{q+r}^{q+r+s} : \mathcal{R}_{q+r+s}(\omega) \rightarrow \mathcal{R}_{q+r}(\omega) \quad \forall r, s \in \mathbb{N}$$

are all surjective.

# Equivalence

## Theorem

Two geometric objects  $\omega$  and  $\gamma$  on  $\mathcal{F}$  are equivalent if:

- All invariants  $\psi$  on  $\mathcal{F}$  coincide for some  $x, y \in X$ :

$$\psi(\omega(y)) = \psi(\gamma(x)),$$

- $\mathcal{R}_q(\omega), \mathcal{R}_q(\gamma)$  are integrable with  $c: \mathcal{F} \rightarrow \mathcal{F}_1$ :

$$(I \circ j_1)(\omega) = c(\omega),$$

$$(I \circ j_1)(\gamma) = c(\gamma).$$

## Example

In the introductory example  $(\omega, \Omega)$  is equivalent to  $(\bar{\omega}, \bar{\Omega})$  if:

$$d\omega = c_1\Omega, \quad d\bar{\omega} = c_1\bar{\Omega}$$

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# Groupoids $\Theta_q$ of Gauge Transformations

- A simple observation for LPDOs:

$$(g^{-1}Lg) u(x) = g^{-1}L(gu(x))$$

$\rightsquigarrow$  Find a groupoid containing only transformations:

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- On  $Y = X \times \mathbb{R}$ , coordinates  $(x, u)$ ,  $(y, v)$  define  $\Theta_q \leq \Pi_q(Y, Y)$ :

$$y = x, \quad u v_u = v.$$

$$\Rightarrow v(x, u) = g(x)u.$$

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- All previous slides remain true if  $\Pi_q$  is restricted to  $\Theta_q$ :

$$0 \longrightarrow \mathcal{R}_q(\omega) \longrightarrow \Theta_q \begin{array}{c} \xrightarrow{\Phi_\omega} \\ \xrightarrow{\omega} \end{array} \mathcal{F}$$

- Natural  $\Theta_q$ -bundles:  $\Theta_q$ -action on  $\mathcal{F}$ .

# Laplace Example – Natural Bundle

- Construct a natural  $\Theta_q$ -bundle for:

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c.$$

- Gauge transformation:

$$\begin{aligned}g^{-1}Lg &= D_{x^1}D_{x^2} + \left(\frac{g_{x^2}}{g} + a\right)D_{x^1} + \left(\frac{g_{x^1}}{g} + b\right)D_{x^2} \\ &\quad + \frac{g_{x^1x^2}}{g} + a\frac{g_{x^1}}{g} + b\frac{g_{x^2}}{g} + c\end{aligned}$$

- Reminder:  $\Theta_q : y = x, \quad u v_u = v.$   
 $\Rightarrow v(x, u) = g(x)u.$
- Natural  $\Theta_q$ -bundle  $\mathcal{F} = Y \times \mathbb{R}^3$ , coordinates  $(x, u, a, b, c)$ :

$$a = \left(\frac{v_{x^2}}{v} + \hat{a}\right),$$

$$b = \left(\frac{v_{x^1}}{v} + \hat{b}\right),$$

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# Laplace Example – Prolongation

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c$$

- $\Theta_q$ -action on  $\mathcal{F} = Y \times \mathbb{R}^3$ :

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- Coordinates of  $J_1(\mathcal{F})$ :

$$a_{x^1}, \quad a_{x^2}, \quad a_u, \quad b_{x^1}, \quad b_{x^2}, \quad b_u, \quad c_{x^1}, \quad c_{x^2}, \quad c_u.$$

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- Coordinates of  $J_{1,X}(\mathcal{F})$ :

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- Coordinates of  $J_{1,X}(\mathcal{F})$ :

$$a_{x^1}, \quad a_{x^2}, \quad b_{x^1}, \quad b_{x^2}, \quad c_{x^1}, \quad c_{x^2}.$$

## Theorem

If the  $\Theta_q$ -action on  $\mathcal{F}$  depends on  $v_{x^\mu}/v$  only then

- Sections  $\omega(x)$  of  $\mathcal{F}$  are well-defined.
- $j_1(\omega)(x)$  restricts to the  $\Theta_q$ -subbundle  $J_{1,X}(\mathcal{F}) \subseteq J_1(\mathcal{F})$ .

# Invariants on $\Theta_q$ -bundles

- $\Theta_q$  is defined by:

$$y = x, \quad u v_u = v.$$

- Invariant  $\psi : J_{r,X}(\mathcal{F}) \rightarrow \mathbb{R}$ .
- Derivatives  $D_{x^i} \psi$  are invariants on  $J_{r+1,X}(\mathcal{F})$ .

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- If all coordinates of  $\mathcal{F}_1$  are invariants,

$$\mathcal{F}_1 \cong \mathcal{F} \times \mathbb{R}^k$$

$\Rightarrow$  the invariants on  $\mathcal{F}_1$  are a generating set.

- Next bundle of integrability conditions  $\mathcal{F}_2$ :

$$\mathcal{F}_2 \cong \mathcal{F}_1 \times \mathbb{R}^{nk}$$

# Laplace Example – Vessiot Structure Equations

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c$$

- $\Theta_q$ -action on  $\mathcal{F} = Y \times \mathbb{R}^3$ :

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- Better coordinates of  $\mathcal{F}_1$ :

$$h = c - a_{x^1} - ab, \quad a_{x^2}, \quad b_{x^1}, \quad k = c - b_{x^2} - ab.$$

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- No equivariant sections  $\mathcal{F} \rightarrow \mathcal{F}_1$ .
- Vessiot structure equations on  $\mathcal{F}_2 \cong \mathcal{F}_1 \times \mathbb{R}^4$ :

$$h_{x^1} = c_1(h, k), \quad h_{x^2} = c_2(h, k), \quad k_{x^1} = c_3(h, k), \quad k_{x^2} = c_4(h, k).$$

$\Rightarrow \{h, k\}$  is a generating set of invariants.

# Example from Shemyakova & Winkler [SW07]

## Example

Third order LPDOs under gauge transformations:

$$(D_x + qD_y)D_{xy} + a_{20}D_{xx} + a_{11}D_{xy} + a_{02}D_{yy} + a_{10}D_x + a_{01}D_y + a_{00}$$

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Equivalence conditions:  $\rightsquigarrow$  second order invariants.

Order	$\Sigma$	Invariants
0	2	$q, I_1 = 2a_{20}q^2 - a_{11}q + 2a_{02}$
1	8	$q_x, q_y, I_x^1, I_y^1$ 4 new invariants
2	15	14 = 6 + 8 old, 1 new
3	21	21 old

Generating set from [SW07] is very compact!

The end.

Thanks!

# Literature



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